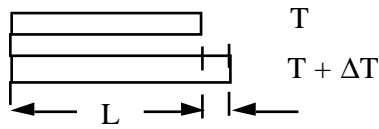


Definition of T:

T is equal in any 2 bodies at thermal equilibrium.

Thermal Expansion



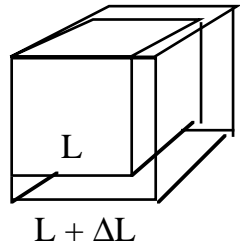
Define

$$\frac{\Delta L}{L} = \alpha \Delta T$$

$$\frac{\Delta V}{V} = \beta \Delta T$$

α is coefficient of linear expansion

β is coefficient of volume expansion



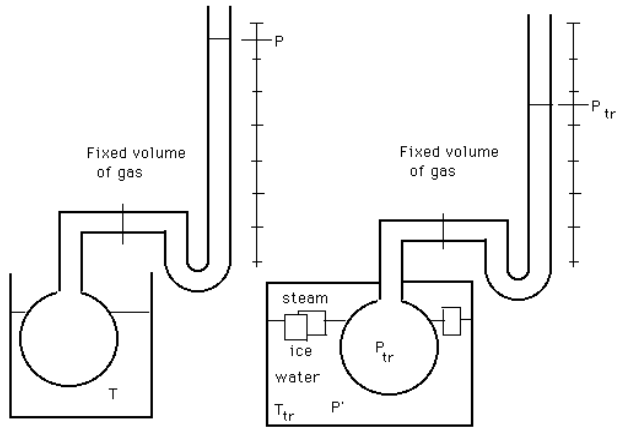
$$\Delta V = (L + \Delta L)^3 - L^3$$

$$= \dots$$

$$\approx V \cdot 3\alpha \Delta T$$

$$\therefore \beta \approx 3\alpha$$

Formal definition of T (Ideal gas thermometer)



$$T = T_{tr} \cdot \lim_{P_{tr} \rightarrow 0} \left(\frac{P}{P_{tr}} \right)_{V}$$

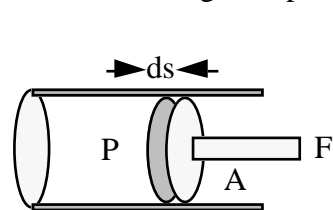
where $T_{tr} = 273.16 \text{ K}$

Heat Capacity: $C = \frac{\Delta Q}{\Delta T}$ extensive quantity (for a body)

Specific Heat: $c = \frac{\Delta Q}{M \Delta T}$ intensive quantity (of a substance)

Latent Heat: heat per unit mass for change of phase (at constant T).

Work done against pressure P



$$dW = \mathbf{F} \cdot d\mathbf{s}$$

$$= PA \, ds$$

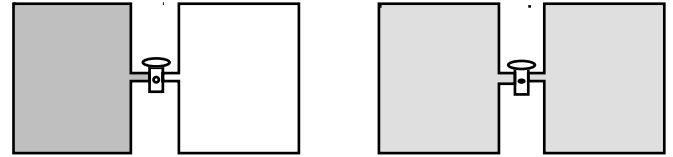
$$= PdV$$

1st Law $dU = dQ - dW$ where U is a state function

Isobaric P const **Isochoric** V constant

Adiabatic Process: $\Delta Q = 0$ (fast or insulated)

Free expansion: open tap



No work done, Experimentally, find $\Delta Q = 0$
 $\therefore \Delta U = 0 \quad \therefore U = U(T)$ only

Calculating work for gases:

Isobaric: $W = \int P \, dV = P \Delta V$

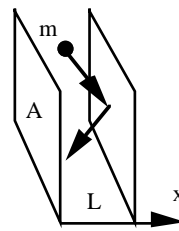
Isothermal: $W = \int P \, dV = \int \frac{nRT}{V} \, dV = nRT \ln \frac{V_f}{V_i}$

Adiabatic: $PV^\gamma = \text{constant}$ revise derivation

$$P = \frac{P_i V_i^\gamma}{V^\gamma} \quad W = \int P \, dV = P_i V_i^\gamma \int \frac{dV}{V^\gamma}$$

$$= (\gamma - 1) P_i V_i^\gamma \left(\frac{1}{V_i^{\gamma-1}} - \frac{1}{V_f^{\gamma-1}} \right)$$

Kinetic Theory



$\Delta \text{momentum} = 2mv_x$
collide every $2L/v_x$

$$|\overline{F}| = \left| \frac{\Delta p}{\Delta t} \right| = \frac{2mv_x}{2L/v_x}$$

$$= \frac{mv_x^2}{L}$$

$$F_{\text{all}} = PA = \frac{Nm \overline{v_x^2}}{L}$$

$v^2 = v_x^2 + v_y^2 + v_z^2$; random motion \Rightarrow

$$\overline{v_x^2} = \frac{1}{3} \overline{v^2}, \text{ so:}$$

$$PAL = PV = \frac{N}{3} m \overline{v^2} \quad P = \frac{1}{3} \rho \overline{v^2}$$

$$\frac{1}{2} m \overline{v^2} = \overline{\epsilon} \text{ and } PV = NkT$$

$$\therefore \overline{\epsilon} = \frac{1}{2} m \overline{v^2} = \frac{3}{2} \frac{PV}{N} = \frac{3}{2} kT$$

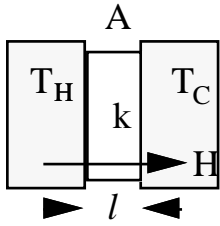
i.e. $\frac{1}{2} kT$ per degree of freedom

$$\frac{1}{2} m \overline{v^2} = \frac{3}{2} kT$$

$$v_{r.m.s.} \equiv \sqrt{\overline{v^2}} = \sqrt{\frac{3kT}{m}}$$

density in field \rightarrow Boltzmann Distⁿ $c = c_0 e^{-E/kT}$

Heat conduction



$$H \equiv kA \frac{T_H - T_C}{l}$$

Planck's radiation law:

$$S(\lambda) = \frac{2\pi c^2 h}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1}$$

$$\frac{\partial S}{\partial \lambda} = 0 \rightarrow \lambda_{\max} = \frac{(2898 \mu\text{m.K})}{T}$$

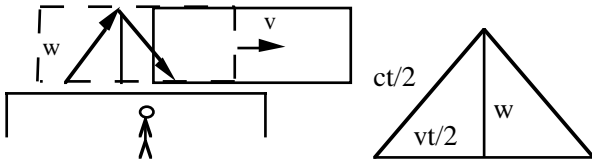
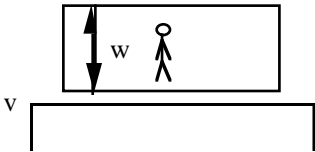
$\int d\lambda \rightarrow \infty$ **Black body radiation**

$$H_{\text{nett}} = \epsilon \sigma A (T^4 - T_{\text{rad}}^4)$$

For both: in **steady state** ($\frac{d}{dt} = 0$), $H_{\text{in}} = H_{\text{out}}$

Principle of special relativity:

Mechanics *and Electromagnetism* are the same in inertial frames.



$$t' = 2w/c \quad t = 2 \frac{\sqrt{w^2 + v^2(t/2)^2}}{c}$$

$$c^2 t^2 = 4 \left(\frac{c^2 t'^2}{4} + \frac{v^2 t^2}{4} \right)$$

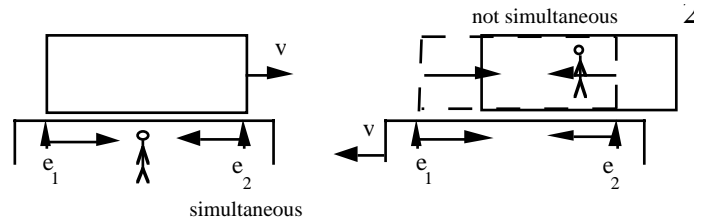
$$\therefore t' = t \sqrt{1 - \frac{v^2}{c^2}}$$

$$t' = t/\gamma \quad (\text{but from the vehicle...})$$

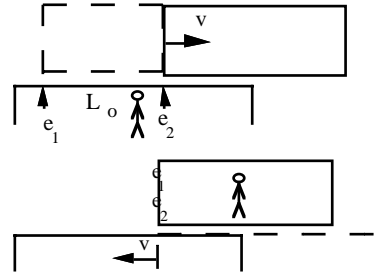
Proper time t_0 in the rest frame.

$$\frac{\Delta t}{\Delta t_0} = \gamma \geq 1 \quad \text{in all other frames } t \text{ is faster than proper time}$$

Relativity of simultaneity and time-order



Proper length L_0 is length in the rest frame



$$L_0 = v \Delta t$$

$$L' = v \Delta t' =$$

$$v \Delta t_0$$

$$\frac{L'}{L_0} = \frac{v \Delta t_0}{v \Delta t} = \frac{1}{\gamma} \leq 1$$

$$x' = Ax + Bt$$

$$y' = y \quad z' = z$$

$$t' = Ct + Dx$$

$$\text{At } x' = 0, \quad x = vt.$$

$$\therefore A = -B/v$$

$$\text{At } t = 0, \quad x' = \gamma x = Ax.$$

$$\therefore x' = \gamma(x - vt)$$

$$\text{At } x = 0, \quad t' = \gamma t$$

$$\therefore C = \gamma$$

$$\text{At } x' = 0, \quad t = \gamma t' \text{ \& } x = vt$$

$$t/\gamma = \gamma t' - Dvt$$

$$\therefore 1 = \gamma^2 - D\gamma v$$

$$\therefore D = \frac{\gamma^2 - 1}{\gamma v} \therefore$$

Lorentz transformations $v \ll c \Rightarrow$ Galileo

$$x' = \gamma(x - vt)$$

$$y' = y \quad z' = z$$

$$t' = \gamma \left(t - \frac{v x}{c^2} \right)$$

$$u_{x'}' = \frac{dx'}{dt'} = \dots = \frac{u_x - v}{1 - \frac{v u_x}{c^2}} \quad \text{not additive}$$

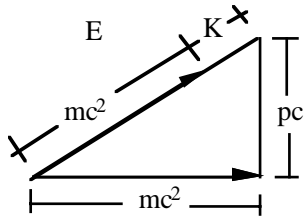
Relativistic Mechanics

Define $p \equiv \gamma m v$. Conserved in both frames.

$$K = \int_{v=0} dW = \dots = mc^2(\gamma - 1)$$

$$\gamma m c^2 = mc^2 + K \quad = \text{proper energy} + K$$

Energy and momentum

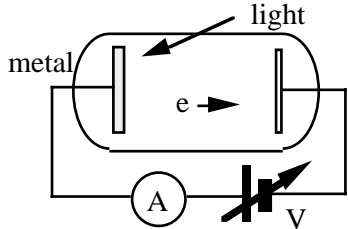


$$p = \gamma mv$$

$$\therefore p^2 c^2 + (mc^2)^2 = \dots = E^2$$

$$E = \sqrt{(pc)^2 + (mc^2)^2}$$

The photoelectric effect



Experiment:
 $eV_0 = hf - hf_0$

$$eV_{\text{stop}} = K \quad hf = hf_0 + K \quad hf_0 = \text{work function}$$

$$\text{Photon energy} \quad E = hf$$

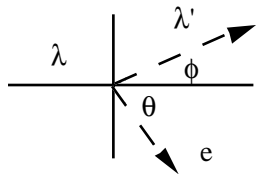
$$p = \frac{E}{c} \quad \text{classical EM}$$

result

for photon $\therefore p = \frac{hf}{c} = \frac{h}{\lambda}$

Compton scattering

X-ray photon from e



$$hf = hf' + mc^2(\gamma - 1)$$

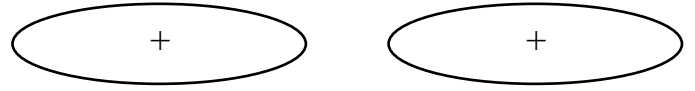
$$\frac{h}{\lambda} = \frac{h}{\lambda'} + mc(\gamma - 1)$$

$$\frac{h}{\lambda} = \frac{h}{\lambda'} \cos \phi + m\gamma v \cos \theta$$

$$0 = \frac{h}{\lambda'} \sin \phi - m\gamma v \sin \theta$$

eliminate $v, \theta \rightarrow \Delta\lambda = \frac{h}{mc} (1 - \cos \phi)$

de Broglie $p = h/\lambda$ for massive particles⁻
Bohr - de Broglie atom



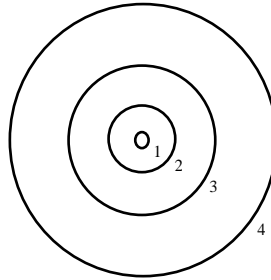
$$2\pi r = n\lambda \quad 2\pi r \neq n\lambda$$

constructive interference destructive interference

Constructive interference, $F_{\text{elec}} = ma = mv^2/r \rightarrow$

$$E = -\left(\frac{me^4}{8\epsilon_0^2 h^2}\right) \frac{1}{n^2}$$

Hydrogen spectrum



Quantum numbers n.

e from n_2 to $n_1 \rightarrow$

$$\frac{1}{\lambda} = \frac{\Delta E}{hc} =$$

$$-R \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

R Rydberg's const

Photon: $\left(\begin{array}{l} \text{probability of} \\ \text{absorption} \end{array} \right) \propto \left(\begin{array}{l} \text{amplitude} \\ \text{of EM field} \end{array} \right)^2$

Massive particles $\left(\begin{array}{l} \text{probability of} \\ \text{'finding' it} \end{array} \right) \propto \Psi\Psi^*$

Heisenberg's Uncertainty Principle

(Finite) Fourier transform: finite length sample \rightarrow limited precision in f. $\delta f \geq 1/2\pi T$

$$E = hf. \quad \therefore \delta E = h\delta f \gtrsim h/2\pi\delta t$$

$$\delta E \cdot \delta t \geq h/2\pi$$

Similarly: $\delta(1/\lambda) \gtrsim 1/2\pi\delta x \rightarrow \delta p \cdot \delta x \geq h/2\pi$

Examples: Heisenberg's microscope, 'Orbit' of the H atom (in ground state), Virtual particles, especially Yukawa's pion, Diffraction through a slit, Young's experiment with localisation

Quantum tunnelling

$$\frac{\hbar^2}{8\pi^2 m} \frac{\partial^2 \Psi}{\partial x^2} + (E - V)\Psi = 0$$

if $K > 0$, Ψ has wave solutions (cos or sin x)

if $K < 0$ Ψ has exponential decay

Examples α decay, tunnel diode, Scanning Tunnelling Electron Microscope