

PHYS1169 Light notes. Joe Wolfe, UNSW

Physical optics (→ interference and diffraction)

Electromagnetic radiation

Speed $c = 3.00 \times 10^8 \text{ ms}^{-1}$

$2.997923458 \times 10^8 \text{ ms}^{-1}$ (This is definition of metre)

λ : for visible light, $400 \text{ nm} < \lambda < 800 \text{ nm}$

$$\text{Typical } f = \frac{c}{\lambda} \sim \frac{3 \times 10^8 \text{ ms}^{-1}}{5 \times 10^{-7} \text{ m}} = 6 \times 10^{14} \text{ Hz} = 600 \text{ THz}$$

See <http://www.phys.unsw.edu.au/~jw/EMSpectrum.html>

Geometrical optics only works if size $\gg \lambda$. Why?

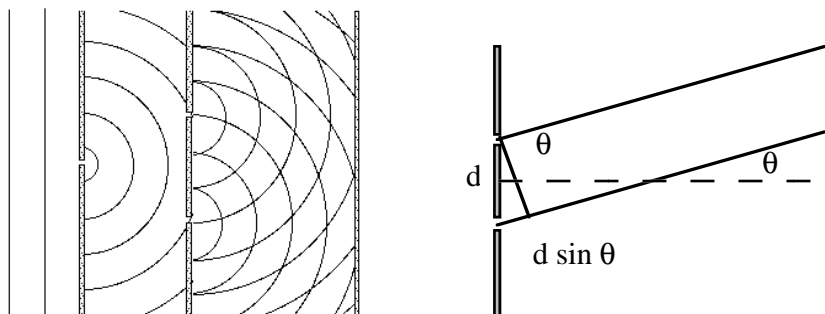
Huygen's principle: each point on a wave front \sim a new source

Is light

- a) a particle (Newton) → geometric optics
- b) a wave (Young) → physical optics
- c) both
- d) none of the above → quantum physics

Hence: **Young's experiment**

Light through one slit (gives *coherent source*) then **two slits** gives interference pattern on screen.



This is the basis of interferometry - powerful tools

Constructive interference if

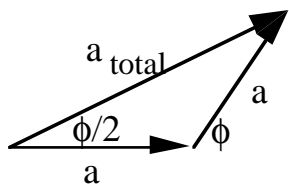
$$d \sin \theta = m \lambda$$

Destructive interference if $d \sin \theta = \left(m + \frac{1}{2}\right)\lambda$

At an angle θ , the phase difference ϕ is

$$\frac{\phi}{2\pi} = \frac{\Delta \text{ pathlength}}{\lambda} = \frac{d \sin \theta}{\lambda}$$

$$\therefore \phi = \frac{2\pi}{\lambda} d \sin \theta \quad \text{how to add two sin waves? phasor diagrams}$$



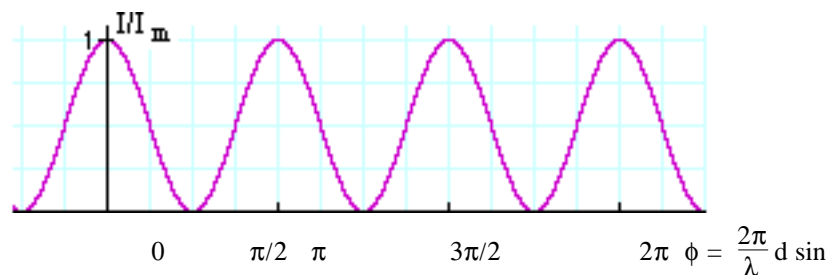
$$a_{\text{tot}} = 2a \cos \beta$$

$$\beta = \phi/2 = \frac{\pi}{\lambda} d \sin \theta$$

$$\text{Intensity} \propto \text{amplitude}^2 \quad \therefore I \propto 4a^2 \cos^2 \beta$$

$$I_{\text{max}} \text{ if } \beta = \phi/2 = 0$$

$$I = I_{\text{max}} \cos^2 \beta \quad \text{where } \beta = \frac{\pi}{\lambda} d \sin \theta$$



θ

Numerical example:

Example 3 narrow slits radiate uniformly and in phase.

Interference pattern on screen. (a) Show phase diagrams for (i) central maximum (ii) 1st order maximum (iii) minima between i and ii.

(b) Sketch $I(\phi)$

i) $\phi = 0$, central max

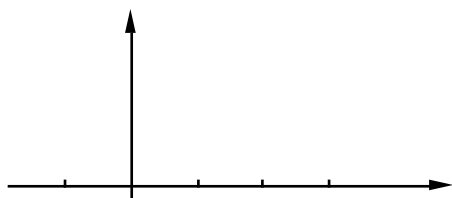
ii) $\phi = 2\pi$, 1st order max

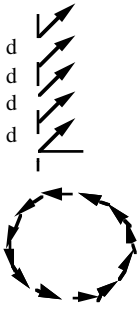
iii) ϕ

$\phi = 2\pi/3 \quad \phi = 4\pi/3 \quad \text{both} \rightarrow 0 \text{ resultant}$

b) if $\phi = \pi$, amplitude = 1/3 central max

$$\therefore \text{intensity} = I_{\text{max}}/9$$





Diffraction grating has very many slits. Used to measure λ very accurately.

If there are N slits per unit length, $d = 1/N$.

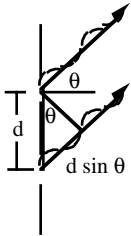
The first minimum is *very* close (small ϕ to close polygon), ie very narrow maxima

For constructive interference

$$d \sin \theta = m\lambda$$

$m = 1 \rightarrow$ 1st order spectrum

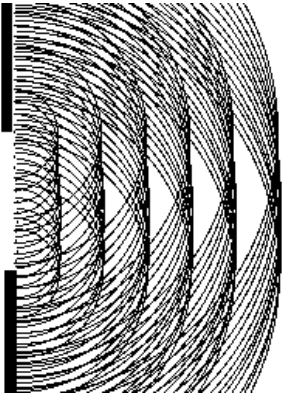
$m = 2 \rightarrow$ 2nd order spectrum



$$\theta_{\text{red}} = \sin^{-1} \frac{m\lambda_{\text{red}}}{d}$$

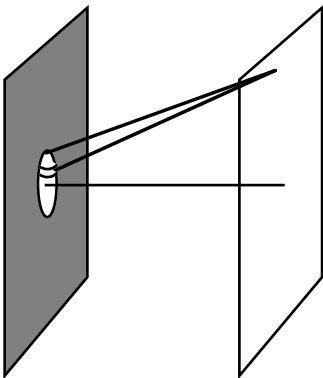
$$\theta_{\text{blue}} = \sin^{-1} \frac{m\lambda_{\text{blue}}}{d}$$

Diffraction from a slit (Later from a circle)

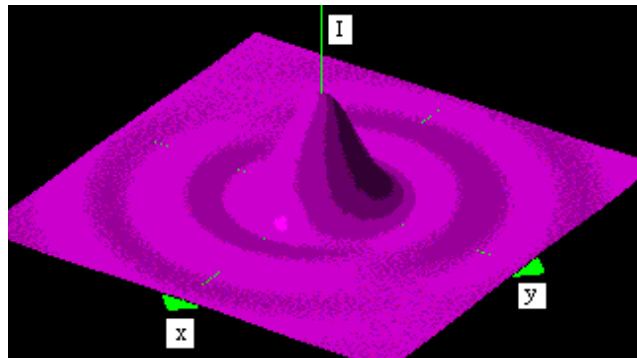


Use Huygen's construction: beam of finite width, interference if $d \sim \lambda$

Circular aperture:

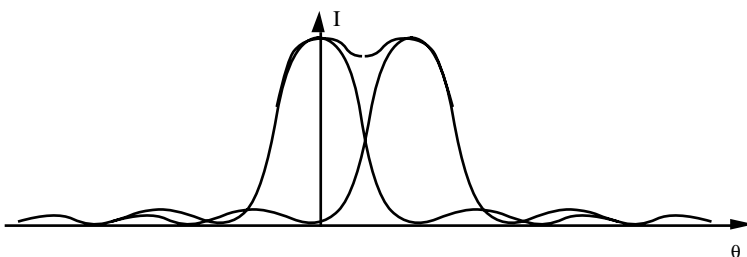


messy integral \rightarrow



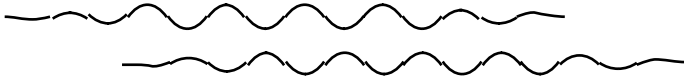
First minimum at $\sin \theta = 1.22 \frac{\lambda}{d}$ d is diameter

can resolve $\sim \theta$ with lens diam d . Rayleigh's criterion



Coherence length

Wave trains have finite length - coherence length



Only interfere if $\Delta \text{ pathlength} < l$

Examples

Radio transmitter $E = E_m \sin(kx - \omega t)$ *E_m, ω, k vary slowly*

Hot body (e.g. lamp) $l \sim 1$ m, but different regions have different, random phase. \therefore , to get interference, use a pin hole and keep $\Delta \text{ path} \ll 1$ m

LASER (Light Amplification by Stimulated Emission of Radiation) $l \gg \text{km}$

Interference in thin films

Light in a medium travels at c/n .

$$\lambda_{\text{medium}} = \lambda/n$$

$$\Delta\phi = 2\pi \frac{\Delta \text{ path}}{\lambda_{\text{med}}} = 2\pi \frac{n \Delta \text{ path}}{\lambda}$$

Define **Optical path length** $\equiv n \cdot \text{pathlength}$

$$\Delta\phi = 2\pi \frac{\Delta \text{ optical pathlength}}{\lambda}$$

Reflections *remember reflections in strings*

From less dense to more dense

$$\Delta\phi = \pi$$

From more dense to less dense

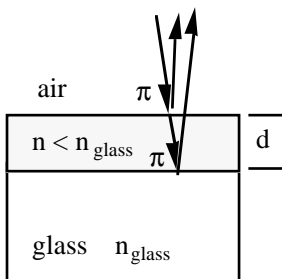
$$\Delta\phi = 0$$

transmitted wave

has no phase change

Example: Non reflective coating

useful in camera lens etc



Coating has $1 < n < n_{\text{glass}}$

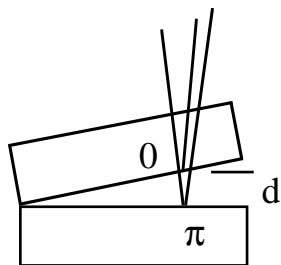
How thick should it be to give minimum reflection?

Both reflections have π phase change. As $d \rightarrow 0$, \rightarrow constructive interference.

For destructive, we want $\Delta \text{ optical pathlength} = \lambda/2$

$$2nd = \lambda/2 \quad d \sim \frac{\lambda}{4n} = \frac{550 \text{ nm}}{4 \cdot 1.2} \sim 115 \text{ nm}$$

Air wedge



Destructive interference if

$$2d = m \lambda$$

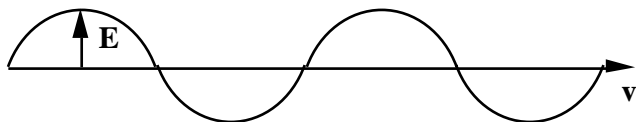
Constructive interference if

$$2d = \left(m + \frac{1}{2}\right)\lambda$$

At L.H. end, destructive (dark), then count fringes to get thickness. *(see corridor display)*

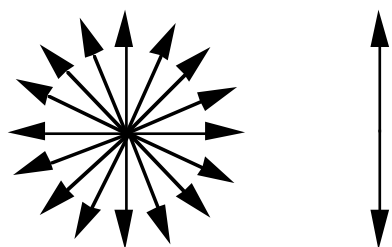
Also: Newton's rings, oil slicks

Polarisation.

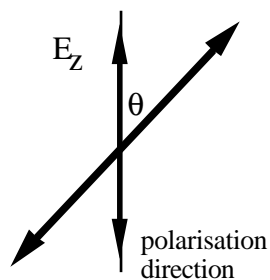


EM waves are **transverse** waves, \therefore can be polarised. Usually light has waves with **E** in all directions

Unpolarised Plane polarised



Polaroid materials allow **E** in only one dirn



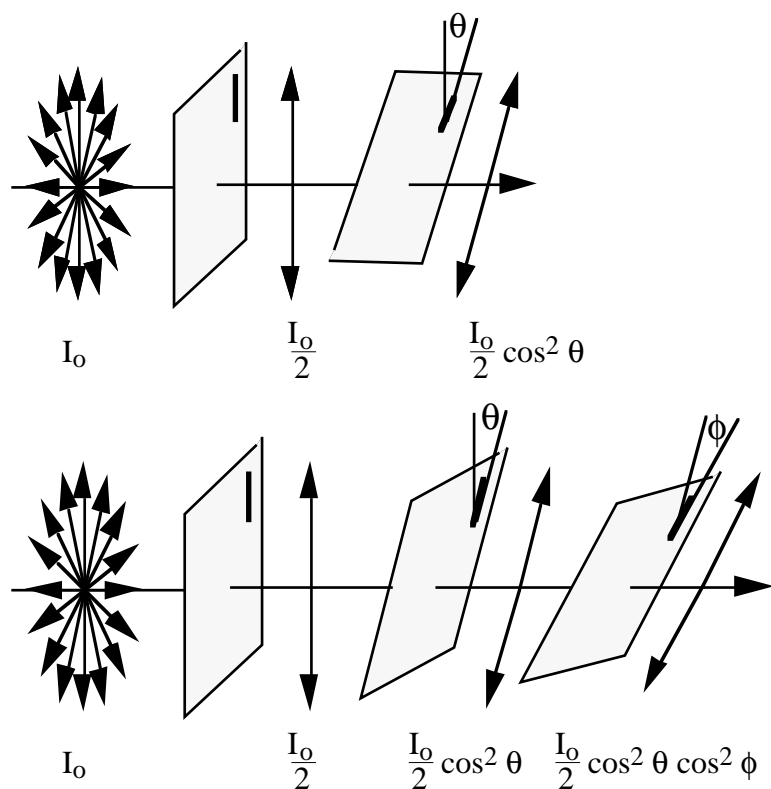
$$E_{\text{transmitted}} = E \cos \theta$$

Malus' Law:

$$I_{\text{trans}} = I_{\text{in}} \cos^2 \theta$$

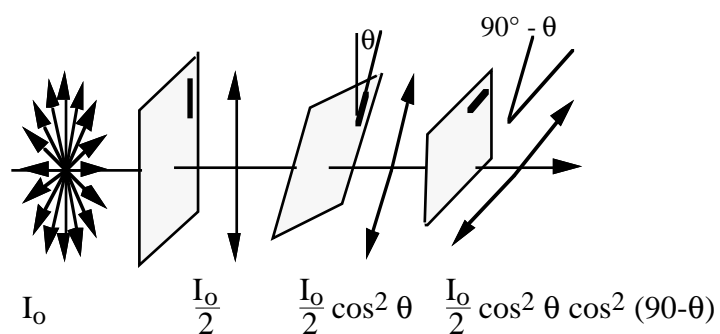
Average of $\cos^2 \theta$ over all angles is $1/2$ \therefore

Unpolarised plane polarised plane polarised



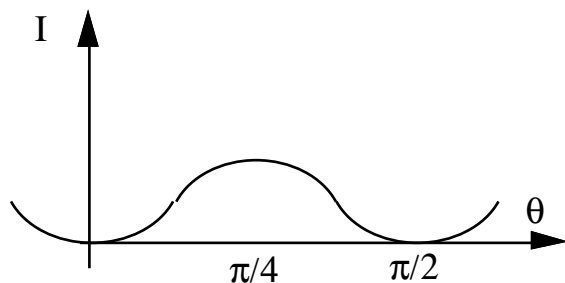
Note: even if $\theta + \phi = 90^\circ$, $I \neq 0$

Example. Two crossed polarisers (polarⁿ dirⁿ are at right angles) have a third polariser between them. What is the transmitted intensity? What angle for the middle polariser gives maximum transmission?



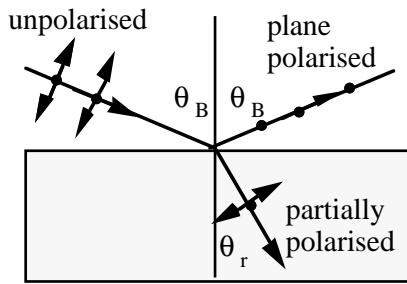
$$I = \frac{I_0}{2} \cos^2 \theta \sin^2 \theta$$

$$= \frac{I_0}{8} \sin^2 2\theta$$



Polarisation by reflection

For wave in medium, \mathbf{E} of light causes oscillation $\parallel \mathbf{E}$. This oscillation can produce (only) transverse waves, hence polarisation of reflected wave.



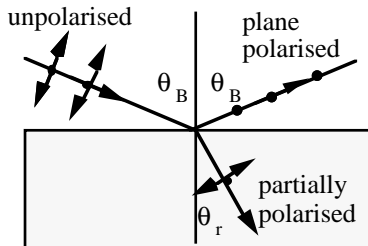
When refracted ray \perp reflected \rightarrow plane polarised reflected wave (Brewster's angle θ_B).

\rightarrow polaroid sunglasses (*see your optics kit*)

Scattering of light also polarises

More effective for short λ , \rightarrow blue sky

Example. What is Brewster's angle for a medium with $n = 1.40$?



Refraction:
$$n = \frac{\sin \theta_B}{\sin \theta_r}$$

If refracted and reflected are at 90° ,

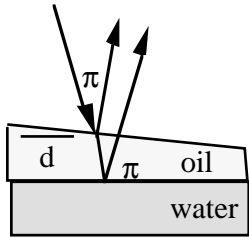
$$\theta_B + \theta_r = 90^\circ$$

so
$$\sin \theta_r = \cos \theta_B$$

$$n = \frac{\sin \theta_B}{\cos \theta_B} = \tan \theta_B$$

$$\theta_B = \tan^{-1} 1.4 = 54^\circ$$

Example. An oil slick ($n = 1.20$) floats on water. What are the thicknesses for which red light ($\lambda \cong 700 \text{ nm}$) is reflected weakly? What does the slick look like at its thinnest point?



$$n_{\text{water}} > n_{\text{oil}}$$

Constructive interference if

$$\Delta \text{OPL} = m \lambda$$

Destructive interference if

$$\Delta \text{OPL} = \left(m + \frac{1}{2}\right) \lambda$$

i) If red has destructive interference,

$$\Delta \text{OPL} = 2nd = \left(m + \frac{1}{2}\right) \lambda_{\text{red}}$$

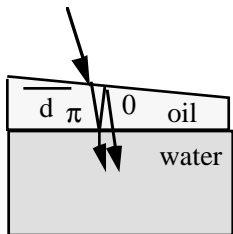
$$d = \frac{\lambda_{\text{red}}}{2n} \left(m + \frac{1}{2}\right)$$

$$m = 0, \quad m = 1, \quad m = 2 \quad \dots$$

$$= 150 \text{ nm}, 440 \text{ nm}, 730 \text{ nm} \quad \text{etc}$$

ii) If $d \ll \lambda$, π phase difference on both paths so constructive interference for all λ , so it looks bright and 'white'.

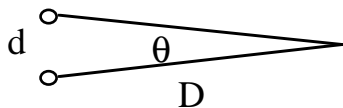
Example. Same problem, but for scuba diver!



Destructive interference if $\Delta \text{OPL} = m \lambda$

Constructive interference if $\Delta \text{OPL} = \left(m + \frac{1}{2}\right) \lambda$

Example. A binary star has an orbital radius of 100 light seconds, and is 10 light years from us. How big must an optical telescope be to resolve them? (take $\lambda = 550 \text{ nm}$). A radio telescope (take $\lambda = 21 \text{ cm}$).



$$100 \text{ s} < 10 \text{ yr} \rightarrow \theta \text{ is small}$$

$$\text{Rayleigh criterion: } \sin \theta_R = 1.22 \frac{\lambda}{a}$$

$$\sin \theta_R \cong \theta_R \cong \frac{100 \text{ light seconds}}{10 \text{ light years}}$$

$$1.22 \frac{\lambda}{a} = \frac{c \cdot 100 \text{ seconds}}{c \cdot 10 \text{ years}}$$

$$\lambda = 550 \text{ nm} \rightarrow a = 2 \text{ m} \quad \lambda = 21 \text{ cm} \rightarrow a = 800 \text{ km}$$

