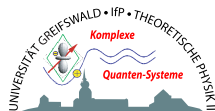


# BOSON-CONTROLLED QUANTUM TRANSPORT



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Topic: Quantum particle strongly interacting with a correlated/fluctuating background medium



$$|\hat{\odot} \cdot \cdot\rangle \rightarrow |* \hat{\odot} \cdot\rangle \rightarrow |* * \hat{\odot}\rangle \rightarrow |* * \hat{\odot}^{\circledast}\rangle \rightarrow |\hat{\odot} * *\rangle \rightarrow |\cdot \hat{\odot} *\rangle \rightarrow |\cdot \cdot \hat{\odot}\rangle$$

in many cases motion bears resemblance to the “Echternacher Springprozession” 😊

related publication: A. Alvermann, D. M. Edwards, HF, Phys. Rev. Lett. **88**, 056602 ('07)

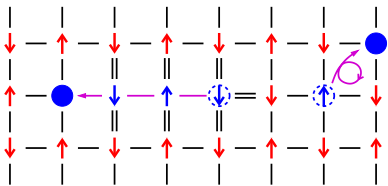
Gordon Godfrey Workshop on Strong Electron Correlations  
UNSW, Sydney, 2007



# MOTIVATION I

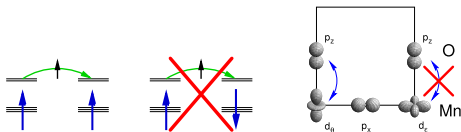
## Strongly correlated charge transport in doped Mott insulators:

### ▶ high- $T_c$ cuprates (AFM spin background)



*classical spins*: “string effect”  
 hole is bound to its starting point  
*quantum spins*: “fluctuations”  
 spin lattice can heal itself with rate  
 controlled by exchange parameter  
 $\leadsto$   $t - J$ -type models

### ▶ colossal magnetoresistive manganites (FM spin background)



strong Hund's rule coupling  
 $\leadsto$  **double-exchange model**

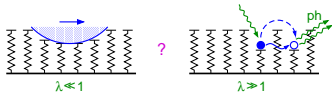
in addition: orbital anisotropy of  
 hopping & EP (JT) coupling

Spin/orbital degrees of freedom might be represented by (e.g. Schwinger) bosons!

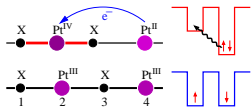


# MOTIVATION II

- Charge transport in systems coupled to phonon or bath degrees of freedom
  - ▶ polarons/excitons, also in CDW materials, DNA, ... (deformable lattice)

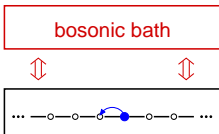


polaron motion - diagonal vs nondiagonal transitions (band vs hopping transport)



~ Holstein-, Fröhlich-, SSH- or Peierls-Hubbard-type models

- ▶ low-D systems – nanowires, quantum dots (disorder, phonons,  $T > 0$ )



system, contacts/leads, bath, ...

appropriate “microscopic” description/modelling?

Again the transport is strongly boson affected, maybe even fluctuation-induced, but now the correlations within the “background” might be weak or even absent!



# MODEL HAMILTONIAN I

How to capture this great variety of transport phenomena in a simplified model?

Let's consider the following rather general (spinless!) Hamiltonian:

$$H = -t_b \sum_{\langle i,j \rangle} c_j^\dagger c_i (b_i^\dagger + b_j) - \lambda \sum_i (b_i^\dagger + b_i) + \omega_0 \sum_i b_i^\dagger b_i + \frac{N\lambda^2}{\omega_0}$$

hopping      boson relaxation      boson energy

Electron emits or absorbs a local boson every time it hops between lattice sites [but hopping creates (destroys) a boson only on the site the particle leaves (enters)!]:

$$\begin{aligned}
 R_i &= c_{i+1}^\dagger c_i b_i^\dagger & | \cdot \hat{\odot} \cdot \rangle & \mapsto | \cdot \star \odot \rangle \\
 L_i &= c_{i-1}^\dagger c_i b_i^\dagger & | \cdot \hat{\odot} \cdot \rangle & \mapsto | \odot \star \cdot \rangle \\
 L_i^\dagger &= c_i^\dagger c_{i-1} b_i & | \hat{\odot} \star \cdot \rangle & \mapsto | \cdot \odot \cdot \rangle \\
 R_i^\dagger &= c_i^\dagger c_{i+1} b_i & | \cdot \star \hat{\odot} \rangle & \mapsto | \cdot \odot \cdot \rangle
 \end{aligned}$$

- $\lambda = 0$  – model is analogous to the classical spin model  $\rightsquigarrow$  “string effect”?
- $\lambda > 0$  allows a boson to decay spontaneously  $\rightsquigarrow$  healing of the “spin lattice”



# MODEL HAMILTONIAN II

- Note that “ $R_i^{(6)} = L_{i+2}^\dagger L_{i+1}^\dagger R_i^\dagger L_{i+2} R_{i+1} R_i$ ” acts as “ $c_{i+2}^\dagger c_i$ ”:

$$|\hat{\odot} \cdot \cdot\rangle \rightarrow |* \hat{\odot} \cdot\rangle \rightarrow |* * \hat{\odot}\rangle \rightarrow |* \hat{\odot}^*\rangle \rightarrow |\hat{\odot} * \cdot\rangle \rightarrow |\cdot \hat{\odot} \cdot\rangle \rightarrow |\cdot \cdot \odot\rangle$$

$$\left| \begin{array}{c} \uparrow \downarrow \\ \odot \uparrow \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \odot \downarrow \\ \uparrow \uparrow \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \downarrow \odot \\ \uparrow \uparrow \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \downarrow \uparrow \\ \uparrow \odot \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \downarrow \uparrow \\ \odot \uparrow \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \odot \uparrow \\ \downarrow \uparrow \end{array} \right\rangle \rightarrow \left| \begin{array}{c} \uparrow \odot \\ \downarrow \uparrow \end{array} \right\rangle$$

$\rightsquigarrow$  lowest order vacuum-restoring process: 1D analogue of 2D “Trugman path”!

- Unitary transformation  $b_i \mapsto b_i + t_f/2t_b$  of  $H$

$$H' = -t_f \sum_{\langle i,j \rangle} c_j^\dagger c_i - t_b \sum_{\langle i,j \rangle} c_j^\dagger c_i (b_i^\dagger + b_j) + \omega_0 \sum_i b_i^\dagger b_i$$

- Different from the  $t$ - $J$  model physics of  $H^{(')}$  is governed by *two* energy ratios:  $t_b/t_f$  and  $t_b/\omega_0$ , where  $t_f = 2\lambda t_b/\omega_0$ !

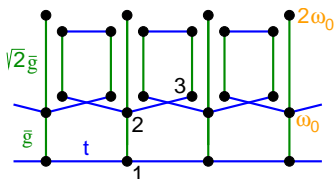
Obviously  $H'$  ( $H$ ) captures the interplay of “coherent” and “incoherent” transport channels realized in many condensed matter systems!



What does it mean: “Solution”?

- Ground state properties

Adapt **variational Hilbert space construction** developed for the Holstein/JT polaron problem (see, e.g., Ku, Trugman, Bonča: PRB 65, 174306 ('02) ):



- |1⟩ e<sup>-</sup> at site 0 with no phonon excitation
- |2⟩ e<sup>-</sup> and phonon at site 0
- |3⟩ e<sup>-</sup> at site 1 and one phonon at site 0

i.e., vertical bonds create or destroy phonons  
act  $m$  times with off-diagonal terms + all translations on an infinite lattice

One- (two-) particle sector: In most cases  $10^4$ - $10^6$  basis states are sufficient to obtain an *8-16 digit accuracy* for  $E_0$ ,  $\langle 0 | \dots | 0 \rangle$ , ... in *any dimension!*  
Note that  $E_0$  calculated this way is *variational* for the *infinite* system!

- Spectral properties at  $T=0$ , thermodynamics, bath degrees of freedom,...

Employ **Kernel Polynomial Method** designed for high-resolution applications: resolution  $\propto 1/\text{number of moments!}$  (cf. review WJAF: RMP 78, 275 ('06))  
**Chebyshev Space Method** – Alvermann, HF: arXiv:0709.3583 [cond-mat.str-el]



# PHYSICAL QUANTITIES OF INTEREST

- ground state energy  $E_0$ , kinetic energy part  $E_{\text{kin}} = \langle 0 | H - \omega_0 \sum_i b_i^\dagger b_i | 0 \rangle$
- quasiparticle band dispersion  $E(\mathbf{k})$ , effective mass  $1/m^* = \frac{\partial^2 E(\mathbf{k})}{\partial k^2} |_{\mathbf{k}=0}$
- particle-boson correlation function  $\chi_{ij} = \langle 0 | b_i^\dagger b_i c_j^\dagger c_j | 0 \rangle$
- one-particle spectral function  $A(\mathbf{k}, \omega) = \sum_n |\langle n | c_{\mathbf{k}}^\dagger | \text{vac} \rangle|^2 \delta[\omega - \omega_n]$
- optical conductivity  $\text{Re}\sigma(\omega) = 2\pi D \delta(\omega) + \sigma_{\text{reg}}(\omega)$ ,

$$\text{regular part } \sigma_{\text{reg}}(\omega) = \pi \sum_{n>0} \frac{|\langle n | j | 0 \rangle|^2}{\omega_n} [\delta(\omega - \omega_n) + \delta(\omega + \omega_n)] ,$$

$$\text{where } j = j_f + j_b \quad \text{with} \quad j_f = it_f \sum_i c_{i+1}^\dagger c_i - c_i^\dagger c_{i+1}$$

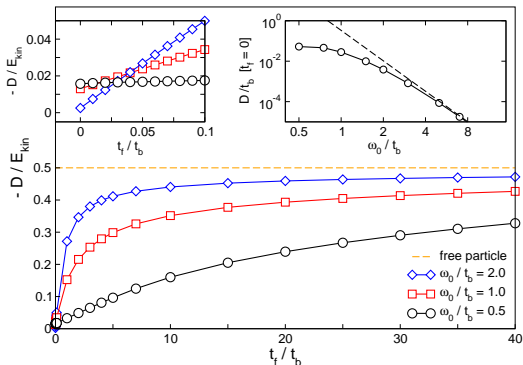
$$j_b = it_b \sum_i c_{i+1}^\dagger c_i b_i^\dagger - c_i^\dagger c_{i+1} b_i - c_{i-1}^\dagger c_i b_i^\dagger + c_i^\dagger c_{i-1} b_i$$

- f-sum rule:  $\int_{-\infty}^{\infty} \sigma(\omega) d\omega = 2\pi D + 2 \int_0^{\infty} \sigma_{\text{reg}}(\omega) d\omega = -\pi E_{\text{kin}}$

$\rightsquigarrow$  consistency check: Drude weight  $D = 1/2m^*$  (Kohn's formula) ✓



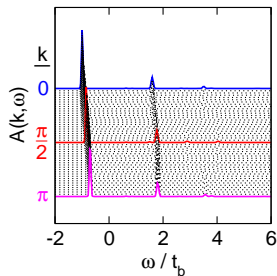
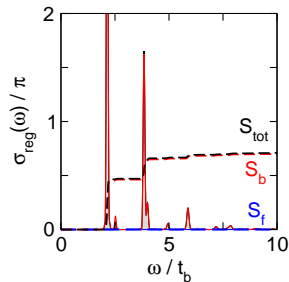
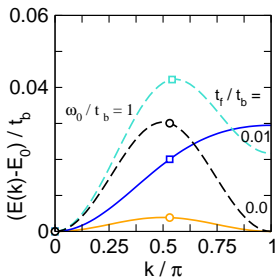
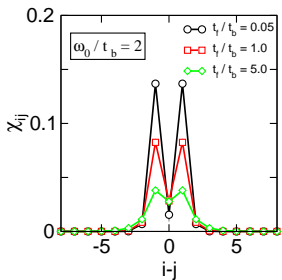
- $D$  scaled to the kinetic energy:



- ▶ free particle:  $t_b = 0 \rightsquigarrow D = t_f$ , i.e.,  $-D/E_{kin} = 0.5$
- ▶ weight of lowest order (vacuum restoring) process scales as  $t_b^6/\omega_0^5 \rightsquigarrow$  boson assisted transport dominates for large  $(t_b/\omega_0)^5(t_b/t_f)$
- ▶  $D$  at  $t_f = 0$  saturates for  $\omega_0 \rightarrow 0$



# CORRELATION-DOMINATED REGIME



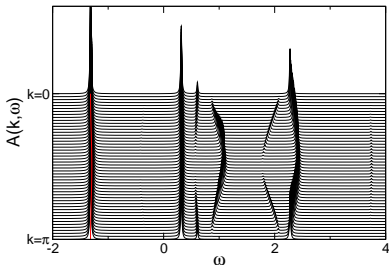
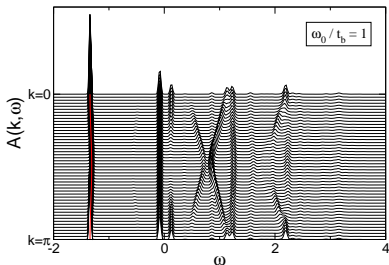
$t_f \leq t_b$  ( $\omega_0$  not too small):

- pronounced NN particle-boson correlations
- strongly renormalised but well-defined quasiparticle band (reminiscent of spin polaron in the t-J model)
- optical response - threshold given by  $\omega_0$
- $\sigma^{\text{reg}} \simeq \sigma_b^{\text{reg}}$   
 $S_{\text{tot}}(\omega) = \int_0^\omega \sigma^{\text{reg}}(\omega') d\omega'$
- $A(k, \omega)$  signals coherent transport

$\rightsquigarrow$  "collective" particle-boson dynamics!



# LIMIT $t_f = 0$ ( $\lambda = 0$ )



## ► exact numerical solution

- particle is still itinerant (but  $D$  is small)
- incoherent contributions
- $k \rightarrow k + \pi$  symmetry

## ► $m$ -boson analytical solution

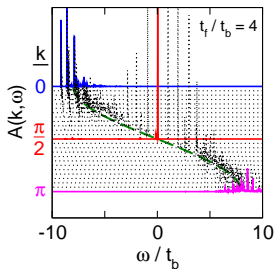
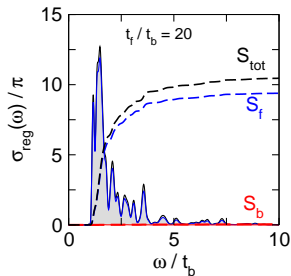
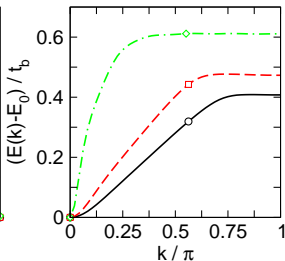
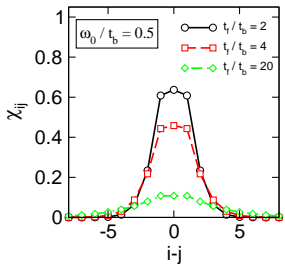
Green function decomposition technique  $\leadsto$  (matrix) continued fraction representation

$m \leq 3$  exact solution possible; only a finite number of states is accessible for the infinite system

$m \geq 4$  infinitely many states will survive



# FLUCTUATION-DOMINATED REGIME



$t_f \gg t_b$  ( $\omega_0$  rather small):

- bosons form a cloud around the particle but are not further correlated
- band flattening near the Brillouin zone boundary

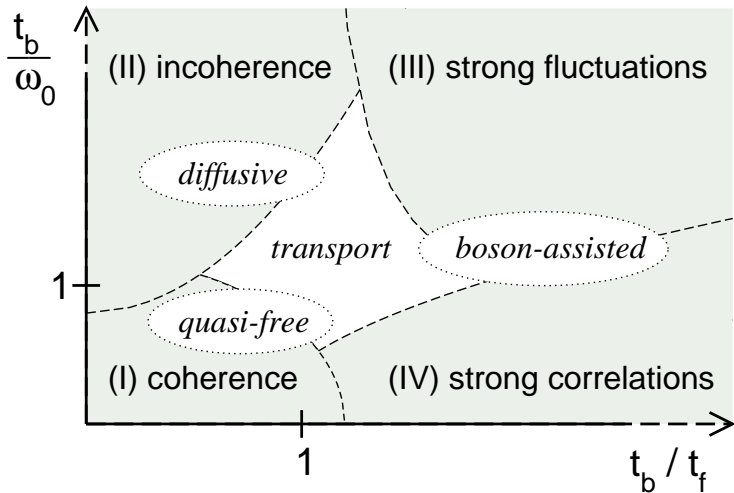
Both is reminiscent of large lattice polarons e.g. in the Holstein model!

- optical response - broad absorption feature
- overdamped character of  $A(k, \omega)$  near  $k = 0, \pi$
- system is almost transparent at  $k = \pi/2$

$\rightsquigarrow$  "diffusive" transport!



# TRANSPORT REGIMES





The model that we propose provides a reduced but realistic description of fundamental aspects of transport in the presence of bosonic fluctuations.

Exact numerical solution ( $N \rightarrow \infty$ )  $\rightsquigarrow$  surprisingly rich physics:

- moving particle creates local distortions of substantial energy in the medium, which may be able to relax
- their relaxation rate determines how fast the particle can move
- “free” particle  $\Leftrightarrow$  magnetic polaron  $\Leftrightarrow$  lattice polaron
- coherent (correlated)  $\Leftrightarrow$  incoherent (diffusive) transport
- bosonic fluctuations act in two competing ways:  
limit transport & assist transport!

And all this is obtained for just one particle! Plus background! 😊

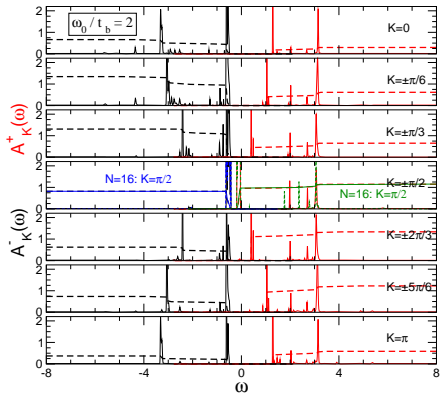
*What about finite carrier density? Half-filled band case? Quantum phase transitions?*





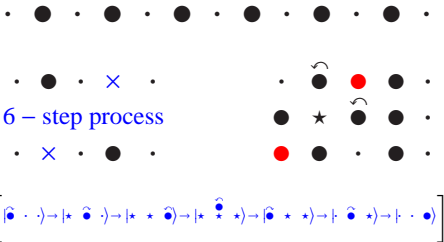
• Doping the insulating state...

$$\lambda = 0.01 \quad (t_f/t_b = 0.01)$$



flat (lower) - **dispersive** (upper) band

dominant hopping processes starting from an AB structure:



► correlation-dominated regime  $\rightsquigarrow$  asymmetric “band structure”

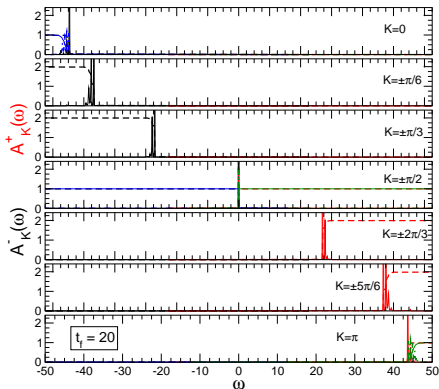


# METAL INSULATOR TRANSITION II

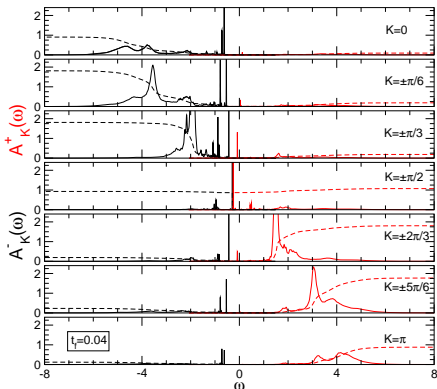
- Fluctuation-dominated regime?

$$\omega_0/t_b = 0.5$$

$\lambda = 5$



$\lambda = 0.01$



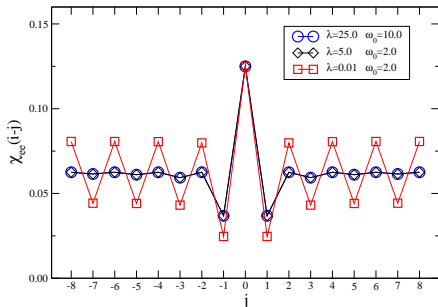
- MIT is suppressed as  $\omega_0 \searrow$  at fixed  $\lambda$  !



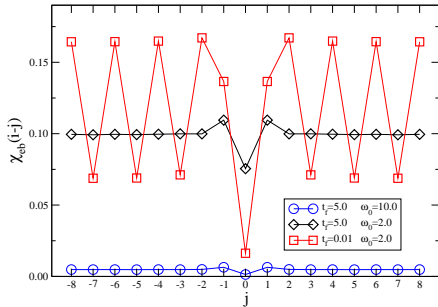
# GROUND-STATE CORRELATIONS

- Particle-particle and particle-boson correlation functions?

$$\chi_{ee}(i-j) = \frac{1}{N_e^2} \sum_i \langle n_i n_{i+j} \rangle$$



$$\chi_{eb}(i-j) = \frac{1}{N_e} \sum_i \langle n_i b_{i+j}^\dagger \rangle$$



► at the MIT long-range order develops !



# SUMMARY – OPEN PROBLEMS

We studied the interplay of collective dynamics and damping in the presence of correlations and fluctuations within a two-channel transport model.

The newly proposed model covers basic aspects of very different Hamiltonians: Hubbard,  $t - J \dots$ , Fröhlich, Holstein,  $\dots$ , SSH - type.

- Influence of spatial dimensionality? VED ✓
- What about two particles - binding? VED ✓
- Finite-density effects ( $0 \leq n \leq 1$ )? DMRG (✓)
- Finite-temperature effects? QMC ?
- Different lattice structures? Frustration? ?
- We need a better analytical understanding of the model, maybe at least for some important limiting cases!

[1D,  $\lambda = 0$  & 3-4 boson approximation,  $\dots$ , semiclassical limit (?),  $\dots$ ]

Everybody is invited to contribute...

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